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On the Information Data Length Theory of Human Emotions

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Abstract

The current paper provides a cutting-edge information data length-theoretic approach to human emotions. This is a revolutionary breakthrough that unifies information data length and the dynamics of emotions. The significant impact of both time and its mean value associated on both the upper and lower bounds of the obtained information data length is observed and noted. The paper concludes with some challenging open problems, combined with concluding remarks and future research pathways.

Keywords: Information Data Length, Psychological Aspects of Human Emotions, Emotional Distribution.

1. Introduction

1.1 Modelling Human Emotions

A mathematical model of human emotion processing in group settings, focusing on how individuals receive and react to emotions from events was developed. It explains how emotions are modelled using a one-dimensional random walk or Wiener process and represented by fixed probability distributions [1]. The author demonstrates that when individuals form a group, the distribution of emotions can also be represented by a fixed distribution, providing insights into emotional responses and behaviours within social dynamics.

Investigated the mathematical modelling of human emotions and decision-making processes within group dynamics [1-7]. It highlights the application of decision field theory to understand how individuals make decisions and how mathematical models are used in psychology and other fields to study dynamic events and decision-making processes. The research aims to provide

insights into human behaviours and emotional responses within social contexts through mathematical analysis and modelling techniques [1]. While mathematical models have been applied in experimental psychology and Bayesian modelling, there seems to be a gap in integrating mathematical methods with the study of human emotions and behaviours. This highlights a potential area for further research and exploration in the field of psychology.

Going through the hypothetical scenario where an event (X) triggers an emotion (x), which then leads to a reaction denoted as F(x) [1]. This relationship is illustrated in Figure 1, showing how events, emotions, and reactions are interconnected. The author suggested that individuals determine whether an event produces positive or negative emotions and adjust their emotional state towards a stable and neutral state (x = 0) through a series of reactions and emotional responses, as illustrated by figure 1(c.f., [1]).

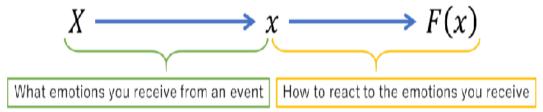


Figure 1: The Relationship Between Events, Emotions, and Reactions in the Context of Individual Emotion Processing

In the context of personal emotion processing, individuals determine whether events elicit positive or negative emotions and whether they lead to large or small emotional or behavioural responses [1]. The process involves randomly starting with an emotion value and then adding 1 or -1 to aim for emotional balance at x = 0, representing stability and emotionlessness. This one-dimensional random walk model suggests that emotions are processed by continuously adding or subtracting values, with the mean emotion value being 0 and variance increasing with the number of processing iterations.

This explains that the process of emotion processing can be likened to a one-dimensional random walk, where positive and negative emotions are determined with a probability of 1/2 and then added together to obtain the total emotion. In this scenario, the mean μ (expected value) of the emotion is 0, indicating a balance between positive and negative emotions, while the variance represents the variability or spread of emotions experienced during the process. This perspective provides a mathematical framework to understand how emotions may fluctuate and balance out over time.

This shows how an individual's reactions are determined by events and emotions they experience, with reactions being like a dictionary once events and emotions are established [1]. It emphasizes that individuals cannot influence events or the type of emotions they receive, but they can influence their reactions based on the initial emotional value assigned to an event. The relationship between the magnitude of emotions and corresponding reactions is highlighted, where strong positive emotions lead to favourable reactions and strong negative emotions lead to unfavourable reactions.

The author had discussed the relationship between a stochastic differential equation and the Fokker-Planck equation in the context of probability density functions and random walk processes, by introducing the concept of a Wiener process, represented by the equation [1]:

$$\chi = \beta W_t \tag{1}$$

where β is a positive real number and W_i is a standard Brownian motion. The equations derived illustrate how changes in probability density functions over time can be described mathematically in the context of random processes.

Thus, the suggested mathematical model can relate to the change in emotions over time, comparing it to a one-dimensional random walk. It introduces equations that describe how emotions change per unit time and how these changes can be represented using partial differential equations. This highlights the relationship between different mathematical expressions and emphasizes the concept of β in the context of the model.

The time-dependent probability distribution of emotions during emotional processing, p(x,t), which provides insights into how emotions evolve over time during the processing phase, helping to understand the average distribution of emotions as the process unfolds. Engaging the mathematical vision provided by, p(x,t) reads as [1]:

$$p(x,t) = \frac{1}{\sqrt{(2\pi t)}} e^{-\left(\frac{x^2}{2t}\right)}$$
 (2)

When individuals are grouped together, the emotion processing of the group is viewed as the average of the individual emotions [1]. This means that the probability density function (PDF) of the group's emotions, $p_{group}(x)$ reads:

$$p_{group}(x) = \frac{1}{\sqrt{(2\pi)}} e^{-\left(\frac{x^2}{2}\right)}$$
 (3)

which describes how emotions are distributed within the group based on statistical principles. This approach simplifies the understanding of how emotions are processed collectively within a group setting, as visualized by figure 2(c.f., [1]).

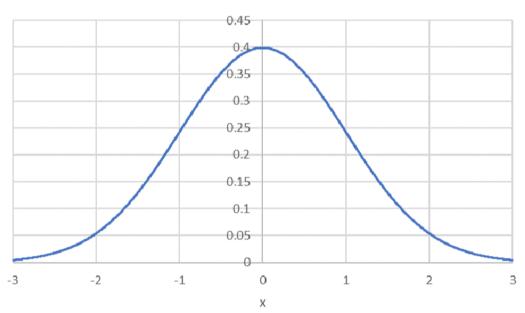


Figure 2: The Portrayed PDF Allows us to Analyse and Compare Emotional Patterns Within Individuals and Groups, Highlighting Similarities in Emotional Processing When People are Grouped Together.

highlighted how individuals and groups share a common equation for expressing emotions, emphasizing the role of understanding social indications in preventing bullying and moral harassment. Some strategies for individuals to adapt to their surroundings, manage their reactions, and handle challenging situations effectively were suggested by [1].

1.2. Information Length Theory

Shannon entropy is not the best descriptor of a time series' statistical variations, which has made using other unique information theoretic notions, including Fisher Information, highly motivating [8-11]. Differential entropy, Kullback-Leibler divergence (KLD), or information length (IL) [12-16]. Time-dependent PDFs provide the ability to trace time series evolution and measure variability, which is the basis of the IL metric's attractiveness [8]. In conclusion, provide an examination of the interval learning (IL) computation of linear stochastic autonomous processes [17–23]. This broadens the scope of application, enabling IL to be used for the abruption of event prediction and facilitating application to different engineering contexts.

1.3. IL as a Concept

Mathematically speaking, if x serves as a nth- order stochastic variable and p(x,t) is a time-dependent PDF of x, then the Information Length L(t) corresponding to its evolution from the initial time $t_0=0$ to the final time $t_F=t$ reads:

$$\mathcal{L}(t) = \int_0^t \frac{dt_1}{\tau(t_1)} = \int_0^t \sqrt{(\varepsilon(t_1))} dt_1 \tag{4}$$

$$\varepsilon(t_1) = \int_{\mathbb{R}^n} \left(\frac{1}{p(x,t_1)} \left[\frac{\partial p(x,t_1)}{\partial t_1} \right]^2 \right) dx \tag{5}$$

provided that $\sqrt{\varepsilon(t)}$ serves as the root-mean-squared fluctuating energy rate.

Having a closer look at (3), it is essential to note that $\tau(t)$ serves as a dynamic temporal unit which provides the correlation time over which the changes of p(x,t) take place [16]. Furthermore, $\tau(t)$ defines the statistical space's time unit. Having said that, $\sqrt{(\varepsilon(t) - 1/(\tau(t)))}$, quantifies the information velocity [18].

It is preferable to compute the underlying value of the mathematical model of the related physical process to comprehend the meaning of \mathcal{L} . Taking the Langevin equation-described first-order stochastic process into consideration:

$$\frac{dx}{dt} = -\gamma(t)(x - f(t)) + \xi \tag{6}$$

x serves as a random variable, f defines a deterministic force, ξ represents a short-correlated random force satisfying that:

$$<\xi(t)\xi(t_1)> = 2D\delta(t-t_1) \text{ and } <\xi(t)> = 0$$
 (7)

where D serves as the amplitude (temperature) of the deterministic force(stochastic nose) ξ .

It is to be noted that equation (7) is so popular to be used as a descriptor of the motion of a particle under a harmonic potential in the form:

$$V(x) = \frac{1}{2}\gamma(x - f(t))^{2}$$
 (8)

Following [23,24], it is found that:

$$\varepsilon(t) = \frac{\left(\frac{d\beta}{dt}\right)^2}{2\beta^2} + 2\beta\left(\frac{dy}{dt}\right)^2 \tag{9}$$

$$p(x,t) = \sqrt{\frac{\beta}{\pi}} e^{-(\beta(x-y))^2}, y(t) = \langle x \rangle = x(0)e^{-G(t)} + F(t)$$

$$\frac{1}{2\beta(t)} = \langle (x(t) - y(t))^2 \rangle = \int_0^t 2D \ e^{-2(G(t) - G(t_1))} dt_1,$$

$$F(t) = \int_0^t e^{-(G(t) - G(t_1))} \gamma(t_1) f(t_1) dt_1, \tag{10}$$

provided that $G(t) = \int_0^t 2D \gamma(t')dt'$.

Considering equation (5), it is clear that $\varepsilon(t)$ is dependent on the changes in both mean and variance defined by the corresponding dynamics of equation (4), portraying p(x,t) three-dimensional space variations, namely (t,x,p(x,t)) as in figure 3 (c.f., [8]),

where the variation of the information velocity, $\sqrt{\varepsilon(t)}$ would occur along the path starting initially from the state probability density function $p(x,t_0)$ to the final state at $p(x,t_F)$ as a descriptor of the speed limit from the statistical deviations of the observables [16].

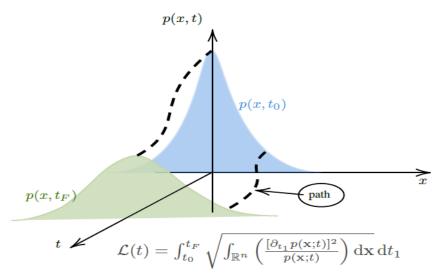


Figure 3

Essentially, we might not be able to see the temporal statistical fluctuations that are occurring when we use differential entropy [12]. This directly follows from the locality's insufficiency because differential entropy primarily measures the differences between any two given PDFs while ignoring any intermediate states [17]. In an alternative symbolism, it merely alerts us to the distinctions that impact the overall evolution of the underlying system. In contrast, any localised changes that take place over the course of the system are measured by IL $\mathcal{L}(t)$, [14,17]. What's more, IL has been hailed as a powerful metric that can bring geometry and stochasticity together, as well as a state-of-theart method for representing an attractor structure [20, 21]. The analogous value of $\lim_{t\to\infty} \mathcal{L}(t)$ would grow linearly from the point of stable equilibrium, based on the location of the mean PDF p(x,0) of the starting state [15,24]. More significantly, using any other information metric results in the loss of this specific trait [17, 23].

Preliminary Theorem(PT)(c.f., [24])

Let f be a function that is defined and differentiable on an open interval (c,d).

If
$$f'(x) > (<0)$$
 for all $x \in (c,d)$, then f is increasing (decreasing) on (c,d) .

The road of this current paper reads: A spotlight introduction is presented in section 1. The key findings are established in section 2. Some challenging open problems, conclusion, and future research work are given in section 3.

2. The Upper and Lower Bounds of the Data Information Length of Human Emotions(HEs), UB_{HEs} , LB_{HEs}

This section is concerned with calculating UB_{HEs} , LB_{HEs} .

2.1 Key Findings

Theorem 1

The IL of HEs , $L_{{\scriptscriptstyle HEs}}$ satisfies the following inequality:

$$t\sqrt{\left(\frac{\lambda^{-3}}{8}(\lambda+1)\right)} > \mathcal{L}_{HES}\left(t,\lambda\right) > t\sqrt{\left(\frac{\lambda^{-3}}{8}\left(1-\lambda^{\frac{3}{2}}\right)\right)}, for some \ \lambda \in (0,t], \lambda \leq 1$$
 (12)

Proof

We have

$$p(x,t) = \frac{1}{\sqrt{(2\pi t)}} e^{-\left(\frac{x^2}{2t}\right)}$$
 (c.f., (2))

Consequently,

$$\frac{\partial p(x,t)}{\partial t} = \frac{t^{-\frac{5}{2}e^{-\left(\frac{x^2}{2t}\right)}}}{2\sqrt{(2\pi)}} (x^2 - t)$$
(13)

$$\frac{\left(\frac{\partial p(x,t)}{\partial t}\right)^{2}}{p(x,t)} = \frac{t^{-\frac{9}{2}}e^{-\left(\frac{x^{2}}{2t}\right)}}{4\sqrt{(2\pi)}}(x^{2}-t)^{2}$$
(14)

It is well known that:

$$\int_0^\infty e^{-u^2} du = \frac{\sqrt{\pi}}{2} \text{ (c.f., [25])}$$
 (15)

Let $I = \int_0^\infty \frac{\left(\frac{\partial p(x,t)}{\partial t}\right)^2}{p(x,t)} dx$. Then,

$$I = \frac{t^{-\frac{9}{2}}}{4\sqrt{(2\pi)}} \left(\int_0^\infty x^2 e^{-\left(\frac{x^2}{2t}\right)} dx + \int_0^\infty t^2 e^{-\left(\frac{x^2}{2t}\right)} dx - 2t \int_0^\infty x e^{-\left(\frac{x^2}{2t}\right)} dx \right)$$
 (16)

Engaging (15), and communicating lengthy steps of mathematical manipulation, yields

$$I = \frac{t^{-\frac{7}{2}}}{4\sqrt{(2\pi)}} \left(\left(\sqrt{\left(\frac{\pi}{2}\right)} \right) t^{\frac{1}{2}} + \left(\sqrt{\left(\frac{\pi}{2}\right)} \right) t^{\frac{3}{2}} - 2t^2 \right)$$
(17)

Communicating(16),

$$I < \frac{t^{-\frac{7}{2}}}{4\sqrt{(2\pi)}} \left(\sqrt{\left(\frac{\pi}{2}\right)}\right) t^{\frac{1}{2}} (t+1) = \frac{t^{-3}}{8} (t+1)$$
(18)

On another separate note,

$$I > \frac{t^{-\frac{7}{2}}}{4\sqrt{(2\pi)}} t^{\frac{1}{2}} \left(\left(\sqrt{\left(\frac{\pi}{2}\right)} \right) t^{\frac{1}{2}} - \left(\sqrt{\left(\frac{\pi}{2}\right)} \right) t^{2} \right) = \frac{t^{-3}}{8} \left(1 - t^{\frac{3}{2}} \right)$$
(19)

Combining (4), (18), and (19), one gets

$$\int_0^t \sqrt{\left(\frac{t^{-3}}{8}(t+1)\right)} dt > \mathcal{L}(t) > \int_0^t \sqrt{\left(\frac{t^{-3}}{8}\left(1 - t^{\frac{3}{2}}\right)\right)} dt$$
 (20)

Following the mean value theorem for definite integrals:

$$t\sqrt{\left(\frac{\lambda^{-3}}{8}(\lambda+1)\right)} > \mathcal{L}_{HES}\left(t,\lambda\right) > t\sqrt{\left(\frac{\lambda^{-3}}{8}\left(1-\lambda^{\frac{3}{2}}\right)\right)} \text{ for some } \lambda \in (0,t]$$
 (c.f., (12)

Notably, $\lambda \le 1$, would guarantee that $LB_{HES} = t\sqrt{\left(\frac{\lambda^{-3}}{8}\left(1-\lambda^{\frac{3}{2}}\right)\right)}$ is real, otherwise $LB_{{\it HEs}}$ would be complex valued, which is not acceptable.

Theorem 2

- LB_{HEs}, UB_{HEs} (c.f., Theorem 1) satisfy: (i) LB_{HEs} is forever decreasing in λ (ii) UB_{HEs} is forever decreasing in λ . (iii) UB_{HEs} is forever increasing in time. (iv) LB_{HEs} is forever increasing in time.

Proof

(i) We have

$$\frac{\partial LB_{HES}}{\partial \lambda} = \frac{t(-3\lambda^{-4} + \frac{3}{2}\lambda^{-\frac{5}{2}})}{4\sqrt{2}\sqrt{\left(\lambda^{-3}\left(1 - \lambda^{\frac{3}{2}}\right)\right)}}$$
(21)

 $\frac{\partial LB_{HES}}{\partial \lambda} < 0 \Leftrightarrow (-3\lambda^{-4} + \frac{3}{2}\lambda^{-\frac{5}{2}}) < 0$, which holds for all $\lambda \in (0,1)$

Communicating (11), and (21), (i) follows.

$$\frac{\partial UB_{HES}}{\partial \lambda} = -\frac{t(3\lambda^{-4} + 2\lambda^{-3})}{4\sqrt{2}\sqrt{(\lambda^{-3} + \lambda^{-2})}}$$
(22)

(ii) Elearly,

 $\frac{\partial UB_{HES}}{\partial \lambda}$ < 0, which implies (ii)

(iii), and (iv) are straightforward.

Define

$$UB_{HES}(\lambda, t) = t\sqrt{\left(\frac{\lambda^{-3}}{8}(\lambda + 1)\right)} \text{ and } LB_{HES}(\lambda, t) = t\sqrt{\left(\frac{\lambda^{-3}}{8}\left(1 - \lambda^{\frac{3}{2}}\right)\right)}$$
 (23)

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2.2. Numerical Experiments

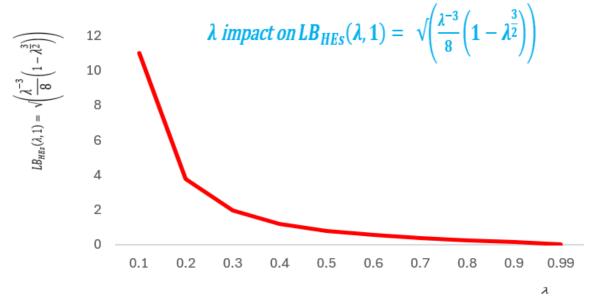


Figure 4

Figure 4 depicts the decreasability of $LB_{HEs}(\lambda,1)=\sqrt{\left(\frac{\lambda^{-3}}{8}\left(1-\lambda^{\frac{3}{2}}\right)\right)}$ in λ , which validates the analytic findings.

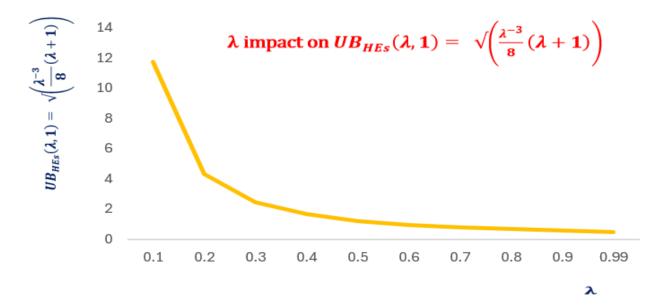
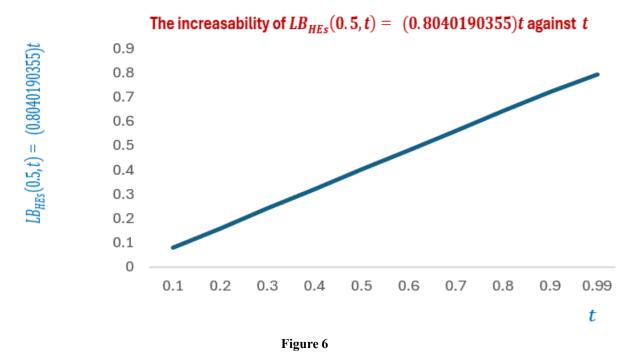


Figure 5

Figure 5 illustrates the decreasability of $UB_{HEs}(\lambda, 1) = \sqrt{\left(\frac{\lambda^{-3}}{8}(\lambda + 1)\right)}$ in λ , which validates the analytic findings.



The numerical findings match the analytic results as shown by figure 6.

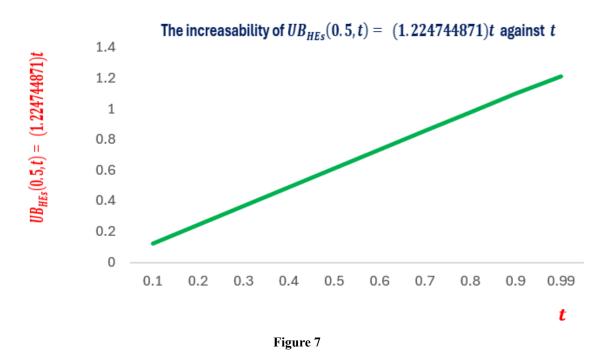


Figure 7 visualizes how $UB_{HEs}(0.5,t) = (1.224744871)t$ increases as times increases.

3. Concluding Remarks, Open Problems, and Future Research

This paper contributes to the establishment of Information Length Theory of HEs. The novel mathematical derivations are undertaken by finding the integral formula of the information length of HEs. Because of the higher complexity to derive the closed form result of the later integral formula, both the upper and the lower bounds of that integral, namely UBHEs, $\textit{LB}_{\textit{HEs}}$ were derived.

Notably, $UB_{HEs}(\lambda,t),LB_{HEs}(\lambda,t)$ are both (λ,t) -dependent. Moreover, these analytic findings were validated numerically.

Here are some emerging open problems:

Open Problem One

Is it mathematically feasible to unlock the challenging problem

of finding the exact analytic form of , $L_{\rm HEs}$ (c.f., (18)), rather than obtaining its upper and lower bounds? This problem is still open.

Open Problem Two

Looking at Eqn (18), can we find any other strict upper and lower bounds for $L_{\rm HFc}$?

Future research pathways include the attempting to solve the proposed open problems and extending the information data length theory to explain other fields of human knowledge, such as engineering, physics and much more.

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